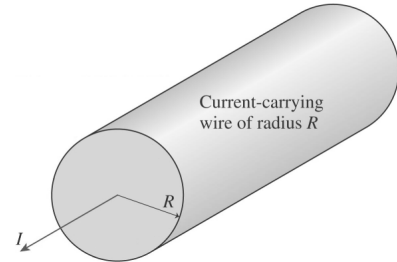


**Example**

An infinitely long current-carrying cable has a radius  $R$  and a current density that varies with distance  $r$  from the central axis according to

$$\vec{J} = J_0 \frac{R}{r} \cos\left(\frac{\pi r}{R}\right) \hat{z}$$

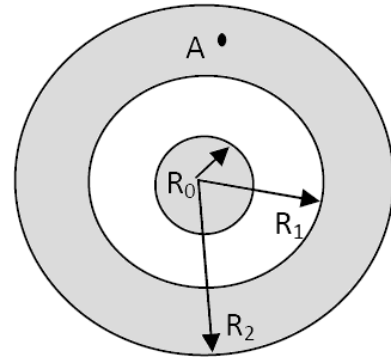


where  $J_0$  is a constant. Find an expression for the magnitude of the magnetic field as a function of  $r$  inside the wire.

**Problem**

A coaxial cable consists of a solid inner wire of radius  $R_0$  and a hollow outer wire with inner radius  $R_1$  and outer radius  $R_2$ , separated by a hollow gap as shown. The inner wire carries a current  $I_0$  out of the page in the positive  $z$  direction. The outer wire has a current density that varies with distance  $r'$  from the central axis according to

$$\vec{J} = -J_0 \frac{R}{r'} \hat{z}$$

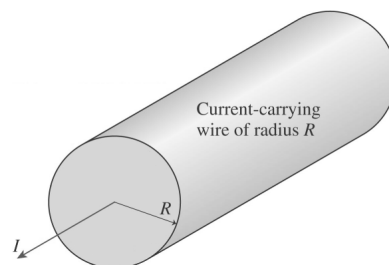


Calculate the magnetic field within the outer cable at point  $A$  (indicted), which is a distance  $R_1 < r < R_2$  from the central axis.

### Example

An infinitely long current-carrying cable has a radius  $R$  and a current density that varies with distance  $r$  from the central axis according to

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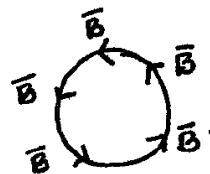
We use Ampere's Law to calculate the magnetic field inside the current cable, say at a distance  $r$  from the center.

Ampere's Law states



that the "circulation of magnetic field along a closed loop is equal to  <sup>$\mu_0$  times</sup> the total current crossing the loop"

$$\Rightarrow \oint_{\text{loop}} \vec{B} \cdot d\vec{L} = \mu_0 I$$



By arguments of symmetry we conclude that the field is along the tangential direction everywhere, equal in magnitude at a given " $r$ " for all  $\phi$  i.e. there would be no angular dependence on the magnitude.

I assume the direction out of the page to be the positive  $\hat{z}$ , thus the loop with this direction is "a counterclockwise" one.



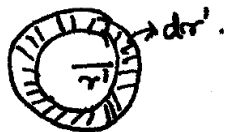
$$\therefore \vec{B}(r) \cdot d\vec{L} = 2\pi r B$$

Now the current enclosed inside this loop of radius

"r" is 
$$\int_{r'=0}^r \vec{J}(r') \cdot d\vec{A}$$

$$d\vec{A} = 2\pi r' dr' \hat{z}$$

(see that it is a vector)



$$\therefore I(r) = \int_0^r J_0 \frac{R}{r'} \cos \frac{\pi r'}{R} dr' \cdot 2\pi r' \hat{z} \cdot \hat{z}$$

$$= J_0 R 2\pi \int_0^r \cos \frac{\pi r'}{R} dr' = 2\pi J_0 R \left. \frac{\sin \frac{\pi r'}{R}}{\pi/R} \right|_0^r$$

$$= 2 J_0 R^2 \sin \frac{\pi r'}{R} \Big|_0^r = 2 J_0 R^2 \sin \frac{\pi r}{R} = I(r)$$

$$\therefore B(r < R) 2\pi r = \left\{ 2 J_0 R^2 \sin \frac{\pi r}{R} \right\} \mu_0$$

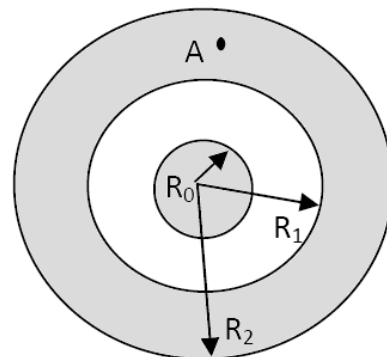
$$\therefore B(r < R) = \frac{\mu_0 J_0 R^2 \sin \frac{\pi r}{R}}{\pi r} \quad B > 0 \Rightarrow \text{in the}$$

positive direction of loop  $\Rightarrow$  +ve  $\hat{\phi}$  direction.

### Problem

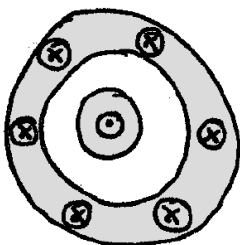
A coaxial cable consists of a solid inner wire of radius  $R_0$  and a hollow outer wire with inner radius  $R_1$  and outer radius  $R_2$ , separated by a hollow gap as shown. The inner wire carries a current  $I_0$  out of the page in the positive  $z$  direction. The outer wire has a current density that varies with distance  $r'$  from the central axis according to

$$\vec{J} = -J_0 \frac{R}{r'} \hat{z}$$



Calculate the magnetic field within the outer cable at point A (indicted), which is a distance  $R_1 < r < R_2$  from the central axis.

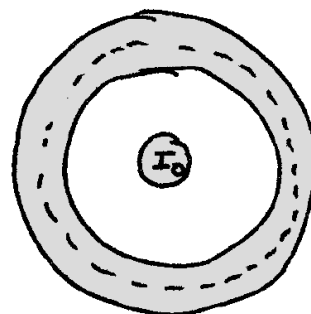
Firstly, we need to realise that the currents in the "core" cable ( $R_0$ ) and the outer cable ( $R_1 < r < R_2$ ) are in opposite directions.



$\otimes \equiv$  into the page ( $-\hat{z}$ )

$\odot \equiv$  out of the page ( $+\hat{z}$ ).

we take a loop in the "counter clockwise" sense. The positive current density direction would then be  $+z$ .



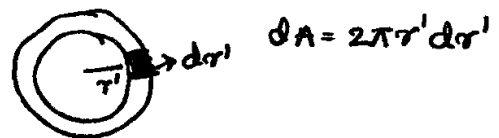
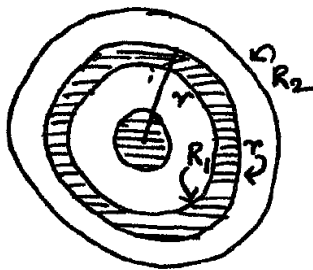
By Arguments of symmetry we again realise that the magnetic field has only a  $\hat{\phi}$  direction and a "radial" dependence i.e  $\vec{B} = B(r) \hat{\phi}$ .

Construct an Amperian loop at a distance "r" within the outer cable  $\Rightarrow r \in [R_1, R_2]$ .

From Ampere's Law we know that

$$\oint_{\text{loop}} \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{enc}} \quad \text{and} \quad \oint \vec{B} \cdot d\vec{l} = B(r) 2\pi r$$

The  $I_{\text{enc}}$  has 2 "contributions". One is due to the current through inner cable and the other is through the outer cable within  $R_1$  &  $r$



$$\begin{aligned} \therefore I_{\text{enc}} &= \underbrace{+ I_0}_{\text{Since along } +z} + \int_{R_1}^r \vec{J} \cdot d\vec{A} \quad \left| \quad d\vec{A} = 2\pi r' dr' \hat{z} \right. \\ &\quad \left. \begin{matrix} \text{(+ve } \hat{z}) \\ \text{(-} \hat{z} \cdot \hat{z} \text{)} \end{matrix} \right. \\ &= I_0 + \int_{R_1}^r J_0 \frac{R_1}{r'} \cdot 2\pi r' dr' (-\hat{z} \cdot \hat{z}) \\ &= I_0 - J_0 R (2\pi) \int_{R_1}^r dr' = I_0 - J_0 2\pi R (r - R_1) \end{aligned}$$

$$\circ \circ \quad B(r) 2\pi r = \left\{ I_0 - J_0 2\pi R(r-R_1) \right\} \mu_0 \quad \#$$

$$B(r) = \frac{I_0 \mu_0}{2\pi r} - \frac{\mu_0 J_0 R(r-R_1)}{r}$$

both the forms of the answer have some very vital information.

①  $B(r)$  could be positive or negative. If it is positive, it means that  $\vec{B}$  is along the positive sense of the loop (counterclockwise). If  $B(r)$  is negative  $\Rightarrow -|B(r)|\hat{\phi}$ , as we can see it is along clockwise direction.

② You can observe the "contribution" to  $\vec{B}$  comes from the inner wire carrying a current  $I_0$  i.e.  $\frac{\mu_0 I_0}{2\pi r}$  as well as from the current in the cable between " $R_1$  &  $r$ ".

③ See that  $R_2$  does not figure in the answer, showing that what matters is the "current enclosed" only.